

DELOCALIZATION OF SCHRÖDINGER EIGENFUNCTIONS

NALINI ANANTHARAMAN

Abstract. A hundred years ago, Einstein wondered about quantization conditions for classically ergodic systems. Although a mathematical description of the spectrum of Schrödinger operators associated to ergodic classical dynamics is still completely missing, a lot of progress has been made on the delocalization of the associated eigenfunctions.

1. Some history

For a long time, the question of the delocalization of eigenfunctions of Schrödinger operators associated to ergodic classical dynamics has been a central problem in mathematical physics. The problem was first posed by Einstein in 1917, who asked whether the eigenfunctions of the Schrödinger operator for a system with ergodic classical dynamics are localized or delocalized. This question was later reformulated in terms of the spectrum of the Schrödinger operator, and it became clear that the delocalization of eigenfunctions is closely related to the spectral properties of the operator.

In the 1970s, the work of von Neumann and others on the spectral theory of operators provided a framework for studying the delocalization of eigenfunctions. The von Neumann algebra associated to a system with ergodic classical dynamics was shown to be a factor, which implies that the eigenfunctions of the Schrödinger operator are delocalized. This result was a major breakthrough in the study of the delocalization of eigenfunctions.

More recently, the work of Anantharaman and others has shown that the delocalization of eigenfunctions is a consequence of the ergodicity of the classical dynamics. This result has been extended to a wide range of systems, including systems with chaotic dynamics and systems with integrable dynamics.

The delocalization of eigenfunctions has important implications for the study of quantum chaos and the spectral properties of Schrödinger operators. It is a topic of active research in mathematical physics, and it is expected that further progress will be made in the near future.

“on the other hand, classical statistical mechanics is essentially only concerned with Type b) [i.e. non integrable systems], for in this case the microcanonical average is the same as the time average”.

Let $H(x; t) = E$, $E \in \mathbb{R}$, (1).

For M

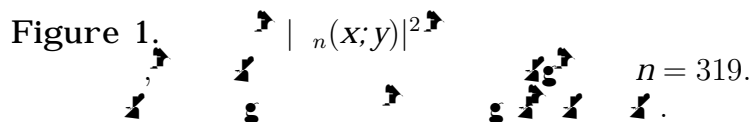
ensembles (fi 2).

\mathbb{R}^2 , chaotic,

103,

arithmetic (“ ” “ ”)

fi 3.



2.1. The role of the geodesic flow . $-\hbar^2 \Delta = E$ (1) ($k = E \hbar^{-2}$) V . $\Delta_k = -k \cdot k$ $\dot{0}$ (-3) $230/$ 11. 2

“ ”.

100 .

$-\hbar^2\Delta$

M

T^*M *cotangent bundle*

$(x; \cdot) \in T^*M$ $x \in M$

$(x; \cdot) \in T^*M$

$t \in \mathbb{R}$, $g^t(x; \cdot) \in T^*M$

$(g^t)_{t \in \mathbb{R}} : T^*M \longrightarrow T^*M$ $g^{t+s} = g^t \circ g^s$ g^0

unit cotangent bundle $S^*M = \{(x; \cdot) \in T^*M; \| \cdot \|_x = 1\}$

g^t .

$(k \longrightarrow +\infty)$,

$\frac{\partial}{\partial t} = i\Delta$

wavefronts

$\epsilon(x=$

$$\begin{aligned}
& - \text{diagram} \quad (2), \quad \text{diagram} \quad \text{ff} \quad \text{diagram} \\
& \quad \quad \quad \sim k^{-1/2} \quad \text{diagram} \quad \text{ff} \quad \text{diagram} \\
& \quad \quad \quad \text{diagram} \quad \text{ff} \quad \text{diagram} \\
& \quad \quad \quad \text{diagram} \quad e^{ict\Delta}, \quad \text{diagram} \quad t_- \\
& \quad \quad \quad \text{diagram} \quad (2) \quad \text{diagram} \quad t_- \\
& \quad \quad \quad \text{diagram} \quad (2) \quad \text{diagram} \quad \text{ff} \quad \text{diagram}
\end{aligned}$$

2.2. L^p -norms as measures of delocalization ?

$$\|k\|_\infty = O(k^{(d-1)/4}):$$
$$\|\lambda\|_{L^p} = O\left(-\frac{(p)}{2}\right)$$

- $(p) = d\left(\frac{1}{2} - \frac{1}{p}\right) - \frac{1}{2}$ for $\frac{2(d+1)}{(d-1)} \leq p \leq +\infty$;
- $(p) = \frac{d-1}{2} \left(\frac{1}{2} - \frac{1}{p}\right)$ for $2 \leq p \leq \frac{2(d+1)}{(d-1)}$.

$$M, \quad \mathbb{S}^d, \quad p_c = \frac{2(d+1)}{(d-1)}, \quad p \geq p_c, \quad p \leq p_c, \quad L^p, \quad p > p_c, \quad L^p, \quad x \in M, \quad \mathcal{L}_x \subset S_x^* M, \quad \mathcal{L}_x = \{v \in S_x^* M; \exists t > 0; g^t(x; v) \in S_x^* M\}, \quad S_x^* M.$$

Theorem 3 ([10] - [10] , [10] , [110]). • Assume there exists a subsequence $n_k \rightarrow +\infty$ and $C > 0$ such that $\|\lambda_{n_k}\|_\infty \geq C^{\frac{(\infty)}{2}}$. Then there exists x such that $x(\mathcal{L}_x) > 0$;

• If M is real analytic, the existence of such subsequence λ_{n_k} is \S to the existence of x such that $\mathcal{L}_x = S_x^* M$, and the first return map $x : S_x^* M \rightarrow S_x^* M$ possesses an absolutely continuous invariant probability measure. (Moreover, in

that case, there exists $t_0 > 0$ such that $g^{t_0}(x; v) \in S_x^*M$ for all $v \in S_x^*M$, that is, there is a common return time).

- If M is real analytic and $\dim M = 2$, the existence of such subsequence λ_{n_k} is due to the existence of $x \in M$ and $t_0 > 0$ such that $g^{t_0}(x; v) = (x; v)$ for all $v \in S_x^*M$.

Anosov property, $\| \lambda \|_{L^\infty} = o\left(\frac{(\infty)}{2}\right) (O(1) o)$.

Theorem 4. (i) If $d = 2$ and M has no conjugate points, or if $d \geq 2$ and M has non-positive sectional curvature, for $p = +\infty$,

$$\| \lambda \|_{L^p} = O\left(\frac{\frac{(p)}{2}}{\sqrt{\log}}\right):$$

(i') Statement (i) actually holds if M has no conjugate points, for all $d \geq 2$.

(ii) (i) holds for all $p > p_c$.

(iii) If M has non-positive sectional curvature, for $p < p_c$, there exists $(p; d) > 0$ such that

$$\| \lambda \|_{L^p} = O\left(\frac{\frac{(p)}{2}}{(\log)^{\sigma(p,d)}}\right)$$

(iv) Statement (iii) still holds for $p = p_c$.

()

2.3. The Shnirelman theorem and the Quantum Unique ergodicity conjecture .

$\int_M |k(x)|^2 d\text{Vol}(x)$
 $k \rightarrow +\infty$, $\int_M |k(x)|^2 d\text{Vol}(x)$

Let $(x_0; 0) \in S^*M$, $a : S^*M \rightarrow \mathbb{R}$, $T \rightarrow +\infty$. Then

$$\lim_{T \rightarrow +\infty} \frac{1}{T} \int_0^T a \circ g^t(x_0; 0) dt = \int_{S^*M} a dL$$

where L is the Liouville measure on S^*M .

Quantum Ergodicity Theorem (Shnirelman theorem).

Theorem 5 ([Shn78, 10, 11]). Let $(M; g)$ be a compact Riemannian manifold, with the metric normalized so that $\text{Vol}(M) = 1$. Call Δ the Laplace-Beltrami operator on M . Assume that the geodesic flow of M is ergodic with respect to the Liouville measure. Let $(\phi_k)_{k \in \mathbb{N}}$ be an orthonormal basis of $L^2(M; g)$ made of eigenfunctions of the Laplacian

$$\Delta \phi_k = -\lambda_k \phi_k; \quad \lambda_k \leq \lambda_{k+1} \rightarrow +\infty.$$

Let a be a continuous function on M . Then

$$(3) \quad \frac{1}{N(\lambda)} \sum_{k, \lambda_k \leq \lambda} \left| \langle \phi_k, a \phi_k \rangle_{L^2(M)} - \int_M a(x) d\text{Vol}(x) \right|^2 \xrightarrow{\lambda \rightarrow +\infty} 0$$

where the normalizing factor is $N(\lambda) = |\{k; \lambda_k \leq \lambda\}|$:

$$\langle \phi_k, a \phi_k \rangle_{L^2(M)} = \int_M a(x) |\phi_k(x)|^2 d\text{Vol}(x).$$

Remark 6. The Cesaro limit (3) implies that there exists a subset $S \subset \mathbb{N}$ of density 1 such that

$$(4) \quad \langle \phi_k, a \phi_k \rangle \xrightarrow{n \rightarrow +\infty, n \in S} \int_M a(x) d\text{Vol}(x).$$

In addition, using the fact that the space of continuous functions is separable, one can actually find $S \subset \mathbb{N}$ of density 1 such that (4) holds for $a \in C^0(M)$. In other words, the sequence of measures $(|\phi_k(x)|^2 d\text{Vol}(x))_{n \in S}$ converges weakly to the uniform measure $d\text{Vol}(x)$.

Actually, the full statement of the theorem says that there exists a subset $S \subset \mathbb{N}$ of density 1 such that

$$(5) \quad \langle \phi_k, A \phi_k \rangle \xrightarrow{n \rightarrow +\infty, n \in S} \int_{S^*M} \sigma^0(A) dL$$

for every pseudodifferential operator A of order 0 on M . On the right-hand side, $\sigma^0(A)$ is the principal symbol of A , that is a function on the unit cotangent bundle S^*M . Equation 4 corresponds to the case where A is the operator of multiplication by the function a .

“ \mathcal{H} ” (\mathcal{H} , \mathcal{H}).
 $\hbar \rightarrow 0$,
 $2, 1, 2$.
 $(\mathcal{H}, \mathcal{H})$,
 $k \rightarrow +\infty$,
 $1, 3$.

Quantum Unique Ergodicity conjecture

Quantum Unique Ergodicity conjecture.

2. . Entropy and support of semiclassical measures.

2. Entropy and support of semiclassical measures. [10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100, 101, 102, 103, 104, 105, 106, 107, 108, 109, 110, 111, 112, 113, 114, 115, 116, 117, 118, 119, 120, 121, 122, 123, 124, 125, 126, 127, 128, 129, 130, 131, 132, 133, 134, 135, 136, 137, 138, 139, 140, 141, 142, 143, 144, 145, 146, 147, 148, 149, 150, 151, 152, 153, 154, 155, 156, 157, 158, 159, 160, 161, 162, 163, 164, 165, 166, 167, 168, 169, 170, 171, 172, 173, 174, 175, 176, 177, 178, 179, 180, 181, 182, 183, 184, 185, 186, 187, 188, 189, 190, 191, 192, 193, 194, 195, 196, 197, 198, 199, 200, 201, 202, 203, 204, 205, 206, 207, 208, 209, 210, 211, 212, 213, 214, 215, 216, 217, 218, 219, 220, 221, 222, 223, 224, 225, 226, 227, 228, 229, 230, 231, 232, 233, 234, 235, 236, 237, 238, 239, 240, 241, 242, 243, 244, 245, 246, 247, 248, 249, 250, 251, 252, 253, 254, 255, 256, 257, 258, 259, 260, 261, 262, 263, 264, 265, 266, 267, 268, 269, 270, 271, 272, 273, 274, 275, 276, 277, 278, 279, 280, 281, 282, 283, 284, 285, 286, 287, 288, 289, 290, 291, 292, 293, 294, 295, 296, 297, 298, 299, 300, 301, 302, 303, 304, 305, 306, 307, 308, 309, 310, 311, 312, 313, 314, 315, 316, 317, 318, 319, 320, 321, 322, 323, 324, 325, 326, 327, 328, 329, 330, 331, 332, 333, 334, 335, 336, 337, 338, 339, 340, 341, 342, 343, 344, 345, 346, 347, 348, 349, 350, 351, 352, 353, 354, 355, 356, 357, 358, 359, 360, 361, 362, 363, 364, 365, 366, 367, 368, 369, 370, 371, 372, 373, 374, 375, 376, 377, 378, 379, 380, 381, 382, 383, 384, 385, 386, 387, 388, 389, 390, 391, 392, 393, 394, 395, 396, 397, 398, 399, 400, 401, 402, 403, 404, 405, 406, 407, 408, 409, 410, 411, 412, 413, 414, 415, 416, 417, 418, 419, 420, 421, 422, 423, 424, 425, 426, 427, 428, 429, 430, 431, 432, 433, 434, 435, 436, 437, 438, 439, 440, 441, 442, 443, 444, 445, 446, 447, 448, 449, 450, 451, 452, 453, 454, 455, 456, 457, 458, 459, 460, 461, 462, 463, 464, 465, 466, 467, 468, 469, 470, 471, 472, 473, 474, 475, 476, 477, 478, 479, 480, 481, 482, 483, 484, 485, 486, 487, 488, 489, 490, 491, 492, 493, 494, 495, 496, 497, 498, 499, 500, 501, 502, 503, 504, 505, 506, 507, 508, 509, 510, 511, 512, 513, 514, 515, 516, 517, 518, 519, 520, 521, 522, 523, 524, 525, 526, 527, 528, 529, 530, 531, 532, 533, 534, 535, 536, 537, 538, 539, 540, 541, 542, 543, 544, 545, 546, 547, 548, 549, 550, 551, 552, 553, 554, 555, 556, 557, 558, 559, 560, 561, 562, 563, 564, 565, 566, 567, 568, 569, 570, 571, 572, 573, 574, 575, 576, 577, 578, 579, 580, 581, 582, 583, 584, 585, 586, 587, 588, 589, 590, 591, 592, 593, 594, 595, 596, 597, 598, 599, 600, 601, 602, 603, 604, 605, 606, 607, 608, 609, 610, 611, 612, 613, 614, 615, 616, 617, 618, 619, 620, 621, 622, 623, 624, 625, 626, 627, 628, 629, 630, 631, 632, 633, 634, 635, 636, 637, 638, 639, 640, 641, 642, 643, 644, 645, 646, 647, 648, 649, 650, 651, 652, 653, 654, 655, 656, 657, 658, 659, 660, 661, 662, 663, 664, 665, 666, 667, 668, 669, 670, 671, 672, 673, 674, 675, 676, 677, 678, 679, 680, 681, 682, 683, 684, 685, 686, 687, 688, 689, 690, 691, 692, 693, 694, 695, 696, 697, 698, 699, 700, 701, 702, 703, 704, 705, 706, 707, 708, 709, 710, 711, 712, 713, 714, 715, 716, 717, 718, 719, 720, 721, 722, 723, 724, 725, 726, 727, 728, 729, 730, 731, 732, 733, 734, 735, 736, 737, 738, 739, 740, 741, 742, 743, 744, 745, 746, 747, 748, 749, 750, 751, 752, 753, 754, 755, 756, 757, 758, 759, 760, 761, 762, 763, 764, 765, 766, 767, 768, 769, 770, 771, 772, 773, 774, 775, 776, 777, 778, 779, 780, 781, 782, 783, 784, 785, 786, 787, 788, 789, 790, 791, 792, 793, 794, 795, 796, 797, 798, 799, 800, 801, 802, 803, 804, 805, 806, 807, 808, 809, 810, 811, 812, 813, 814, 815, 816, 817, 818, 819, 820, 821, 822, 823, 824, 825, 826, 827, 828, 829, 830, 831, 832, 833, 834, 835, 836, 837, 838, 839, 840, 841, 842, 843, 844

“ ”

n_k .

(n_k) .

fi

fi

$g_{\sharp}^t =$, $t \in \mathbb{R}$.

$\rightarrow +\infty$,

2.1.

Theorem 7. 3 Assume M is a compact Riemannian manifold with negative sectional curvature. Assume $\langle n_k; A_{n_k} \rangle$ converges to $\int_{S^*M} {}^0(A) d$ for all A . Then has positive entropy.

,

$h_{KS}(\cdot)$,

$h_{KS}(\cdot) = 0$

$\mapsto h_{KS}(\cdot) \text{ ffi } h_{KS}(\cdot_1 + (1 - \cdot_2) \cdot) = h_{KS}(\cdot_1) + (1 - \cdot) h_{KS}(\cdot_2)$,

$\cdot_1, \cdot_2 \in [0; 1]$

(\cdot) ffi

(\cdot)

$h_{KS}(\cdot) \leq \int \left(\sum_{j=1}^r \cdot_j^+(\cdot) \right) d(\cdot)$

$\cdot_j^+(\cdot)$

$L(\cdot_1, \cdot_2) =$

$d(S^*M, 2d - 1)$, $d - 1$

(\cdot)

$h_{KS}(\cdot) \leq d - 1$;

L .

Corollary 1. Let $\Gamma \subset S^*M$ be the union of all points lying on a periodic trajectory on the geodesic flow (recall, if M has negative curvature, there are countably many periodic geodesics). Let be as in Theorem 7. Then $(\Gamma) < 1$.

, \cdot

$\text{scar}(\cdot, \cdot)$.

strongly scarred

Let Γ be a partial scar, $(\Gamma) > 0$. Then $\text{supp}(\Gamma)$ has Hausdorff dimension > 1 .

Corollary 2. *The support of Γ has Hausdorff dimension > 1 .*

Let Γ be a partial scar, $(\Gamma) > 0$. Then $\text{supp}(\Gamma)$ has Hausdorff dimension > 1 .

Theorem 8 (see [1]). Assume M is a compact Riemannian manifold of dimension d , with constant sectional curvature -1 . Then $\text{supp}(\Gamma)$ has entropy greater than $\frac{d-1}{2}$.

Let Γ be a partial scar, $(\Gamma) > 0$. Then $\text{supp}(\Gamma)$ has Hausdorff dimension > 1 .

Corollary 3. *The support of Γ has Hausdorff dimension $\geq d$.*

Let Γ be a partial scar, $(\Gamma) > 0$. Then $\text{supp}(\Gamma)$ has Hausdorff dimension $\geq d$.

Corollary 4. Let $\Gamma \subset S^*M$ be the union of all points lying on a closed trajectory on the geodesic flow. Let Γ be as in Theorem 7, with M of constant negative curvature. Then $(\Gamma) \leq 1=2$.

Let Γ be a partial scar, $(\Gamma) > 0$. Then $\text{supp}(\Gamma)$ has Hausdorff dimension > 1 .

Theorem 9 (see [1]). $\text{supp}(\Gamma)$ has full support, that is, $(\Omega) > 0$ for any non-empty open set $\Omega \subset S^*M$.

[illegible]

3. Large scale delocalization

[illegible]

3. \dots regular \dots

3.1. Overview of the problem.

$G = (V; E)$.
localized, delocalized?

$I \subset \mathbb{R}$,

- spectral localization :
- exponential localization :
- dynamical localization :

- spectral delocalization :
- ballistic transport :

(G_N)
 (G_N)
 $N \rightarrow \infty$.

$\sum_{x=1}^N |j(x)|^2$
 $1=N$
 M_N
 $N \times N$,
 $(j)_{j=1}^N$

- ∞ norms :

$\|j\|_\infty$

- p norms:

non-ergodic (multi-fractal) $\|j\|_p$

- Scarring :

$(\sum_{x \in \Lambda} |j(x)|^2 \geq)$

- Quantum ergodicity :

$\sum_x a(x) |j(x)|^2$

$a : \{1; \dots; N\} \rightarrow \mathbb{C}$,

most j . all j , quantum unique ergodicity.
 almost sure random random
 2, 3, 33, 1, 1, 13, 1 .1.
 3 , ff

3.2. Entropy.

$(q+1)$ -
 $(q+1)$ -
 $(G_N)_{N \in \mathbb{N}} = (V_N; E_N)$.

$$\mathcal{A}_N f(x) = \sum_{x \sim y} f(y)$$

$$x \sim y \iff x \text{ and } y$$

$$\Delta_N f(x) = \sum_{x \sim y} (f(y) - f(x)) :$$

$$(10) \quad \mathcal{A}_N - (q+1)I = \Delta_N :$$

Theorem 10 (). Let (G_N) be a sequence of $(q+1)$ -regular graphs (with q fixed), $G_N = (V_N; E_N)$ with $V_N = \{1; \dots; N\}$. Assume that¹ there exists $c > 0$; $\epsilon > 0$ such that, for any $k \leq c \ln N$, for any pair of vertices $x, y \in V_N$,

$$(11) \quad |\{\text{paths of length } k \text{ in } G_N \text{ from } x \text{ to } y\}| \leq q^{k(\frac{1-\epsilon}{2})} :$$

Fix $\epsilon > 0$. Then, if f is an eigenfunction of the discrete Laplacian on G_N and if $\Lambda \subset V_N$ is a set such that

$$\sum_{x \in \Lambda} |f(x)|^2 \geq \sum_{x \in V_N} |f(x)|^2 ;$$

then $|\Lambda| \geq N^\alpha$ — where $\alpha > 0$ is given as an explicit function of ϵ ; and c .

$$H_N(f) = -\frac{1}{\log N} \sum_x |f(x)|^2 \ln |f(x)|^2$$

¹This assumption holds in particular if the injectivity radius is $\geq c \ln N$. The interest of the weaker assumption is that it holds for typical random regular graphs [90].

$$\|L^\infty\|_{\infty} = O((\log N)^{-1/4}). \quad (11)$$

3.3. QE on regular graphs. Let $G_N = (V_N; E_N)$ be a $(q+1)$ -regular graph on N vertices.

$$(EXP) \quad \frac{1}{N} \sum_{x \in V_N} \left(\frac{1}{(q+1)^{-1} \mathcal{A}_N} \right)^2 (V_N) \rightarrow 0 \quad \text{as } N \rightarrow \infty.$$

$$\frac{1}{N} \sum_{x \in V_N} \left(\frac{1}{\mathcal{A}_N} \right)^2 (V_N) \rightarrow 0 \quad \text{as } N \rightarrow \infty.$$

(BST) R ,

$$\frac{|\{x \in V_N; (x) < R\}|}{N} \xrightarrow{N \rightarrow \infty} 0$$

$$\frac{(x)}{B(x; r)} \rightarrow 0 \quad \text{as } r \rightarrow \infty.$$

$$(BST) \quad \frac{1}{N} \sum_{x \in V_N} \left(\frac{1}{(q+1)^{-1} \mathcal{A}_N} \right)^2 (V_N) \rightarrow 0 \quad \text{as } N \rightarrow \infty.$$

$$\frac{1}{N} \sum_{x \in V_N} \left(\frac{1}{\mathcal{A}_N} \right)^2 (V_N) \rightarrow 0 \quad \text{as } N \rightarrow \infty.$$

$$(12) \quad \int \left(\frac{1}{N} \sum_{x \in V_N} \left(\frac{1}{\mathcal{A}_N} \right)^2 (V_N) \right) d\mu(x) = \langle \mathcal{A}_x^k \rangle_{\ell^2(x)}.$$

Theorem 11 (). Let $(G_N) = (V_N; E_N)$ be a sequence of $(q+1)$ -regular graphs with $|V_N| = N$. Assume that (G_N) satisfies **(BST)** and **(EXP)**.

Let $(\binom{N}{1}, \dots, \binom{N}{N})$ be an orthonormal basis of eigenfunctions of \mathcal{A}_N in $\mathcal{L}^2(V_N)$.

Let $\mathbf{a}_N : V_N \rightarrow \mathbb{C}$ be a sequence of functions such that $\sup_N \sup_{x \in V_N} |\mathbf{a}_N(\mathbf{x})| \leq 1$: Define $\langle \mathbf{a}_N \rangle = \frac{1}{N} \sum_{x \in V_N} \mathbf{a}_N(\mathbf{x})$.

Then

$$\begin{aligned}
& \frac{1}{N} \sum_{j=1}^N \left| \langle \psi_j^{(N)}; a_N \psi_j^{(N)} \rangle_{\ell^2(V_N)} - \langle a_N \rangle \right|^2 \xrightarrow{N \rightarrow +\infty} 0; \\
(13) \quad & \frac{1}{N} \left| \left\{ j \in [1; N]; \left| \langle \psi_j^{(N)}; a_N \psi_j^{(N)} \rangle_{\ell^2(V_N)} - \langle a_N \rangle \right| > \right\} \right| \xrightarrow{N \rightarrow +\infty} 0; \\
& \langle \psi_j^{(N)}; a_N \psi_j^{(N)} \rangle_{\ell^2(V_N)} = \sum_{x \in V_N} a_N(x) |\psi_j^{(N)}(x)|^2. \\
& \sum_{x \in V_N} |\psi_j^{(N)}(x)|^2 = \frac{1}{N} \sum_{x \in V_N} 1 = 1.
\end{aligned}$$

3. . Non-regular graphs : from spectral to spatial delocalization.

regular

quantum ergodicity theorem

spectral delocalization

spatial delocalization

“If a large finite system is close (in the Benjamini-Schramm topology) to an infinite system having purely absolutely continuous spectrum in an interval I , then the eigenfunctions (with eigenvalues lying in I) of the finite system satisfy quantum ergodicity.”

$$(G_N)_{N \in \mathbb{N}}, \quad V_N, E_N, \quad N \longrightarrow \infty, \quad \text{fix } \rho, \quad G_N, \quad d, \quad \mathcal{A}_N : \mathbb{C}^{V_N} \rightarrow \mathbb{C}^{V_N} \quad |V_N| =$$

$$(\mathcal{A}_N f)(v) = \sum_{w \sim v} f(w);$$

$$V \sim W \quad \text{if} \quad \exists \text{ } \mathcal{A}_N \text{ such that } \left(\frac{\mathcal{A}_N}{N} \right) \rightarrow 0 \text{ as } N \rightarrow +\infty.$$

$$\begin{aligned}
& \text{(EXP)} \quad \mathbb{M} \quad (G_N) \quad \text{fin} \quad (G_N) \\
& \quad \quad \quad P_N : \mathbb{C}^{V_N} \rightarrow \mathbb{C}^{V_N} \\
(1) \quad & (P_N f)(x) = \frac{1}{d_N(x)} \sum_{y \sim x} f(y); \\
& \quad \quad d_N(x) \quad P_N \quad P_N \quad [-1 + \frac{1}{N}; 1 - \frac{1}{N}] \cup \{1\}, \quad > 0 \quad P_N \quad N. \\
& \quad \quad 1 \quad G_N \quad P_N \quad 0 \quad (\quad , \quad 3). \quad (G_N) \\
& \text{(BST)} \quad r > 0, \\
& \quad \quad \lim_{N \rightarrow \infty} \frac{|\{x \in V_N : G_N(x) < r\}|}{N} = 0; \\
& \quad \quad G_N(x) \quad \text{injectivity radius} \quad x, \quad B_{G_N}(x; r) \\
& \quad \quad \text{rooted graphs } (G_N), \quad (\quad 1), \\
& \quad \quad \text{(BST)} \quad \text{tree } [\mathcal{T}; o]. \\
& \text{(BSCT)} \quad (G_N) \quad \mathbb{P} \quad \{[\mathcal{T}; o]\} \quad \mathbb{P} \\
& \quad \quad (\quad)_{j=1}^N \quad \mathcal{A}_N \quad {}^2(V_N). \quad \text{(BSCT)} \quad : \mathbb{R} \rightarrow \mathbb{R}, \\
(1) \quad & \frac{1}{N} \sum_{j=1}^N (\quad)_{j=1}^N \xrightarrow{N \rightarrow +\infty} \mathbb{E}(\langle \cdot, o; (\mathcal{A}_{\mathcal{T}}) \cdot o \rangle) =: \int (t) d\mathbf{m}(t); \\
& \quad \quad \mathcal{A}_{\mathcal{T}} \quad \mathbb{P}. \quad \mathcal{T}, \quad \mathbb{P}. \quad \mathbf{m} \quad \text{integrated density of states} \\
& \quad \quad \mathbb{P}. \quad \mathbb{P}. \quad I \quad \mathcal{A}_{\mathcal{T}} \quad \mathbb{P}. \quad \mathbb{P}. \\
& \quad \quad \mathbb{P}. \quad x, y \in \mathcal{T}, \quad \in \mathbb{C} \setminus \mathbb{R}, \quad (\quad) \\
& \quad \quad \mathcal{G}^\gamma(x; y) = \langle x; (\mathcal{A}_{\mathcal{T}} - \gamma)^{-1} y \rangle_{\ell^2(\mathcal{T})} :
\end{aligned}$$

$$\begin{aligned}
& \mathcal{T} \ni v; w \in \mathcal{T} \quad v \sim w, \quad \mathcal{T}^{(v|w)} \quad \mathcal{A}_{\mathcal{T}^{(v|w)}} \\
& \quad \mathcal{G}^{(v|w)}(\cdot; \cdot; \cdot) \quad \gamma_w(v) := -\mathcal{G}^{(v|w)}(v; v; \cdot). \\
& \text{(Green)} \quad I, \quad s > 0 \\
& \sup_{\lambda \in I, \eta_0 \in (0,1)} \mathbb{E} \left(\sum_{y: y \sim o} |\text{Im} \, o^{\lambda + i\eta_0}(y)|^{-s} \right) < \infty : \\
& \quad \text{(Green)}, \quad \text{fi} \quad (\quad) \\
& [\mathcal{T}; o] \\
& (1) \quad o(J) = \langle o; \mathbb{1}_J(\mathcal{A}_{\mathcal{T}}) o \rangle \quad J \subseteq \mathbb{R} : \\
& \quad \text{(Green)} \quad \sup_{\lambda \in I, \eta_0 > 0} \mathbb{E}(|\mathcal{G}^\gamma(o; o)|^2) < \infty. \\
& \quad \mathbb{P} \dots [\mathcal{T}; o], \quad \frac{1}{\pi} \text{Im} \mathcal{G}^{\lambda + i0}(o; o), \quad \text{(Green)} \quad \mathbb{P} \dots \\
& \quad \mathcal{A}_{\mathcal{T}} \quad I, \quad 12 \quad z \quad \text{(Green)}, \quad s < 0, \\
& \quad 2 \quad \text{(Green)}, \quad \text{fi} \quad I_1 \quad \text{fi} \\
& \quad \bar{I}_1 \subset I. \quad G_N \quad \Gamma_N \setminus \widetilde{G_N} \quad \widetilde{G_N} \\
& \quad (\quad) \quad \tilde{x}; \tilde{y} \quad \tilde{\mathcal{A}}_N \quad \widetilde{G_N}, \quad \in \mathbb{C} \setminus \mathbb{R}, \\
& (1) \quad \tilde{g}_N^\gamma(\tilde{x}; \tilde{y}) = \langle \tilde{x}; (\tilde{\mathcal{A}}_N - \quad)^{-1} \tilde{y} \rangle_{\ell^2(\widetilde{G_N})} :
\end{aligned}$$

Theorem 12 (). Assume that $(G_N; W_N)$ satisfies **(BSCT)**, **(EXP)** and **(Green)**.

Call $(\lambda_j^{(N)})_{j=1}^N$ the eigenvalues of \mathcal{A}_N on $\ell^2(V_N)$, and let $(\gamma_j^{(N)})_{j=1}^N$ be a corresponding orthonormal eigenbasis.

For each N , let $a = a_N$ be a function on V_N with $\sup_N \sup_{x \in V_N} |a_N(x)| \leq 1$.

Then

$$\lim_{\eta_0 \downarrow 0} \lim_{N \rightarrow +\infty} \frac{1}{N} \sum_{\lambda_j^{(N)} \in I_1} \left| \sum_{x \in V_N} a(x) |\gamma_j^{(N)}(x)|^2 - \sum_{x \in V_N} a(x) \frac{N}{\lambda_j^{(N)} + i\eta_0} \gamma_j^{(N)}(x) \right| = 0;$$

for some family of probability measures γ_N on V_N , indexed by a parameter $\gamma \in \mathbb{C} \setminus \mathbb{R}$, defined as follows :

$$\gamma_N(x) = \frac{\text{Im} \, \tilde{g}_N^\gamma(\tilde{x}; \tilde{x})}{\sum_{y \in V_N} \text{Im} \, \tilde{g}_N^\gamma(\tilde{y}; \tilde{y})}.$$

$$(1) \quad \frac{1}{N} \left| \left\{ j^{(N)} \in I_1 : \left| \sum_{x \in V_N} a(x) |j^{(N)}(x)|^2 - \sum_{x \in V_N} a(x) \frac{N}{\lambda_j^{(N)} + i\eta_0} (x) \right| > \right\} \right| \xrightarrow{N \rightarrow +\infty, \eta_0 \downarrow 0} 0 :$$

Corollary 13. *For any $\alpha \in (0;1)$, there exists $c_\alpha > 0$ such that for any $\Lambda_N \subset V_N$ with $|\Lambda_N| \geq N$, we have*

$$(20) \quad \frac{1}{N} \sum_{x \in V_N} F \left(N^{\frac{N}{\lambda + i\eta_0}}(x) \right) \xrightarrow{N \rightarrow +\infty} \mathbb{E} \left(F \left(\frac{\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o)}{\mathbb{E}(\text{Im } \mathcal{G}^{\lambda + i\eta_0}(o; o))} \right) \right) :$$

Remark 14. *The results proven in actually hold for more general Schrödinger operators than adjacency matrices : one can consider weighted Laplacians (with conductances on the edges) and add a potential; in other words, on each G_N , we can consider a discrete Schrödinger operator \mathcal{H}_N . The limiting object in assumption **(BSCT)** is now a random*

rooted tree $[\mathcal{T}; o]$ endowed with a random Schrödinger operator \mathcal{H} . Assumption **(Green)** has to be modified, replacing the adjacency matrix \mathcal{A} by the operator \mathcal{H} . Similarly, in the statement of the theorem, the Green functions \tilde{g}_N to be considered are those of \mathcal{H}_N lifted to the universal cover \widetilde{G}_N .

Remark 15. In particular, our result applies to the case where the limiting system $([\mathcal{T}; o]; \mathcal{H})$ is $\mathcal{T} = \mathfrak{X}$ (the $(q+1)$ -regular tree) with an arbitrary origin o , and $\mathcal{H} = \mathcal{H}_\epsilon = \mathcal{A} + \mathcal{W}$ where \mathcal{W} is a random real-valued potential on \mathfrak{X} . More precisely the values $\mathcal{W}(x)$ ($x \in \mathfrak{X}$) are i.i.d. random variables of common law μ . This is known as the Anderson model on \mathfrak{X} . It was shown by A. Klein [Kl80] that the spectrum of \mathcal{H}_ϵ is a.s. purely absolutely continuous on $I = (-2\sqrt{q} + \epsilon; 2\sqrt{q} - \epsilon)$, provided ϵ is small enough (depending on μ). This just assumes a second moment on μ . Under stronger regularity assumptions on μ , one can show that Assumption **(Green)** holds on I (see [11], following Aizenman-Warzel [2]). Examples of sequences of expander regular graphs G_N with discrete Schrödinger operators \mathcal{H}_N converging to $([\mathfrak{X}; o]; \mathcal{H}_\epsilon)$ are given in [10].

Remark 16. Examples of sequences of (G_N) satisfying our three assumptions were investigated in [11]. In the examples considered there, the limiting trees \mathcal{T} are \mathfrak{X} or \mathfrak{Y} ; roughly speaking, those are trees where the local geometry can only take a finite number of values. If \mathcal{A} is the adjacency matrix of such a tree, we showed in [11] that the spectrum of \mathcal{A} is a finite union of closed intervals, and that there are a finite number of points y_1, \dots, y_ℓ in \mathbb{R} such that Assumption **(Green)** holds on any I of the form $\mathbb{R} \setminus ([y_1 - \epsilon; y_1 + \epsilon] \cup \dots \cup [y_\ell - \epsilon; y_\ell + \epsilon])$ (for any $\epsilon > 0$). We showed – extending Remark 15 – that on such trees, Assumption **(Green)** remains true after adding a small random potential to $\mathcal{A}_\mathcal{T}$. Finally, we showed the existence of sequences (G_N) converging to \mathcal{T} and satisfying the **(EXP)** condition.

Perspectives and link with other work

1. Random regular graphs. [11] deterministic (G_N) , a. “ ”, “ ”, a. random, random, fi, 1, 13, 13, $\{1; \dots; N\}$, $(q+1)$ -, $(q+1)N$, 1, 13, 1, “ ”, adja-
 cency matrix $a_N : \{1; \dots; N\} \rightarrow \mathbb{R}$,

$$\mathcal{N}(0; \frac{2}{j}) \quad 0 \leq j \leq 1. \quad \sqrt{N} \quad j(x) \quad x \quad \{1; \dots; N\} \quad R \geq 0, \quad (\sqrt{N} \quad j(y))_{y, d(y, x) \leq R} \quad G_N, \quad B_{\mathfrak{X}}(0; R).$$

2. From graph Laplacians to Hecke operators.

11. From graph Laplacians to Hodge operators

Let $(g_1, \dots, g_k) \in \Delta_{\mathbb{S}^2}$ be a free subgroup of $SO(3)$.

$$T_k f(x) = \sum_{j=1}^k (f(g_j x) + f(g_j^{-1} x))$$

$$\Delta_{\mathbb{S}^2}.$$

Theorem 19 (Anantharaman, 2010). Assume that g_1, \dots, g_k generate a free subgroup of $SO(3)$.

For each ℓ , let $(\psi_j^{(\ell)})_{j=1}^{2\ell+1}$ be an orthonormal family of eigenfunctions of $-\Delta_{\mathbb{S}^2}$ of eigenvalue $\ell(\ell+1)$, that are also eigenfunctions of T_k .

Then for any continuous function a on \mathbb{S}^2 , we have

$$\frac{1}{2\ell+1} \sum_{j=1}^{2\ell+1} \left| \int_M a(x) |\psi_j^{(\ell)}(x)|^2 d\text{Vol}(x) - \int_M a(x) d\text{Vol}(x) \right|^2 \xrightarrow{\ell \rightarrow \infty} 0.$$

11. T_k is not a pseudodifferential operator, so the argument sketched above to show that the basis $(Y_\ell^m)_{\ell \geq 0, |m| \leq \ell}$ could not satisfy quantum ergodicity does not apply here.

Remark 20. T_k is not a pseudodifferential operator, so the argument sketched above to show that the basis $(Y_\ell^m)_{\ell \geq 0, |m| \leq \ell}$ could not satisfy quantum ergodicity does not apply here.

Remark 21. We note that for very special choices of rotations – rotations that correspond to norm n elements in an order in a quaternion division algebra, the operators T_k are called Hecke operators. It has been conjectured by Böcherer, Sarnak, and Schulze-Pillot [2] that such joint eigenfunctions satisfy the much stronger quantum unique ergodicity property. This conjecture is still open.

3. [1] 2000, (n) infinite, (\cdot, \cdot) almost-every ergodic component, (\cdot, \cdot) M , 10 one

11

3. Quantum ergodicity on Riemann surfaces of high genus.

11 12

11

Theorem 22 ($M \geq 0$). Let (S_N) be a sequence of hyperbolic surfaces, whose genus (equivalently, volume) goes to ∞ .

(EXP) Assume the first eigenvalue $\lambda_1(N)$ of $-\Delta$ on S_N is bounded away from 0 as $N \rightarrow \infty$.

(BSH) Assume there are few short geodesics; in other words, (S_N) converges in the Benjamini-Schramm sense to the hyperbolic disc: for any $R > 0$,

$$\lim_{N \rightarrow +\infty} \frac{\text{Vol}\{x \in S_N : r(x) < R\}}{\text{Vol}(S_N)} = 0$$

where $r(x)$ means the injectivity radius at x .

Fix an interval $I \subset (0; +\infty)$.

Let $(\phi_i^{(N)})$ be an orthonormal basis of eigenfunctions of the Laplacian on S_N .

Let $a = a_N : S_N \rightarrow \mathbb{C}$ be such that $|a(x)| \leq 1$ for all $x \in S_N$. Then

$$\lim_{N \rightarrow +\infty} \frac{1}{\text{Vol}(S_N)} \sum_{\lambda_i(N) \in I} \left| \int_{S_N} a(x) |\phi_i^{(N)}(x)|^2 dx - \langle a \rangle \right|^2 = 0$$

where $\langle a \rangle = \frac{1}{\text{Vol}(S_N)} \int_{S_N} a(x) dx$.

(1; +\infty) L^2- S_N

22

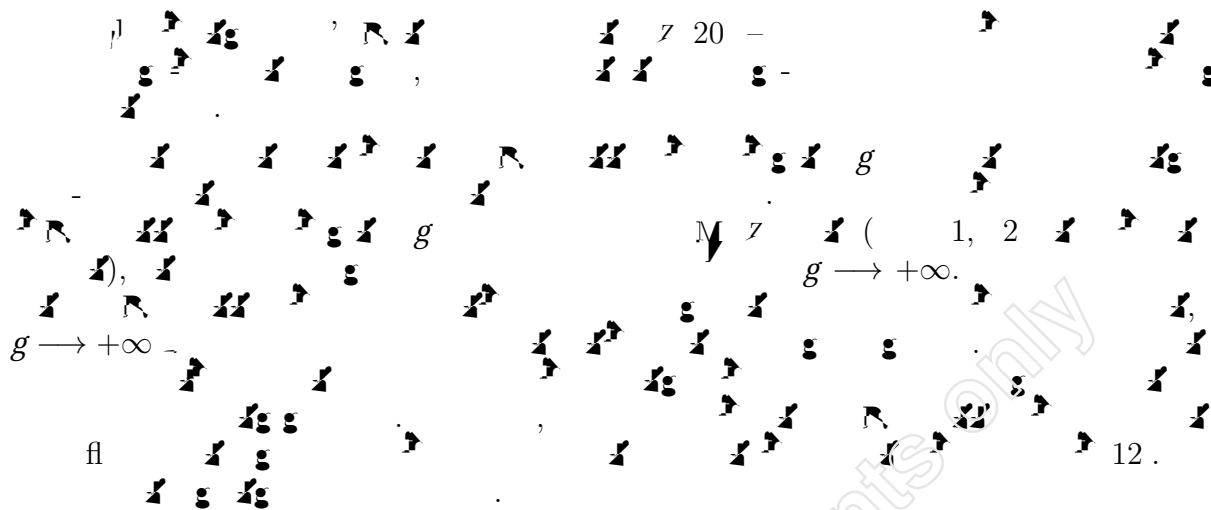
random

random

1

(\phi_i^{(N)})

12



References

- [1] De Luca A., Altshuler B., Kravtsov V. E., and Scardicchio A. Anderson localization on the bethe lattice: Nonergodicity of extended states. *Phys. Rev. Lett.*, 113, 2014.
- [2] Michael Aizenman and Simone Warzel. Absolutely continuous spectrum implies ballistic transport for quantum particles in a random potential on tree graphs. *J. Math. Phys.*, 53(9):095205, 15, 2012.
- [3] Nalini Anantharaman. Entropy and the localization of eigenfunctions. *Ann. of Math. (2)*, 168(2):435–475, 2008.
- [4] Nalini Anantharaman. Quantum ergodicity on regular graphs. *Comm. Math. Phys.*, 353(2):633–690, 2017.
- [5] Nalini Anantharaman and Etienne Le Masson. Quantum ergodicity on large regular graphs. *Duke Math. J.*, 164(4):723–765, 2015.
- [6] Nalini Anantharaman and Stéphane Nonnenmacher. Entropy of semiclassical measures of the Walsh-quantized baker’s map. *Ann. Henri Poincaré*, 8(1):37–74, 2007.
- [7] Nalini Anantharaman and Stéphane Nonnenmacher. Half-delocalization of eigenfunctions for the Laplacian on an Anosov manifold. *Ann. Inst. Fourier (Grenoble)*, 57(7):2465–2523, 2007. Festival Yves Colin de Verdière.
- [8] Nalini Anantharaman and Mostafa Sabri. Poisson kernel expansions for schrödinger operators on trees. *preprint*, <https://arxiv.org/abs/1610.05907>, 2016.
- [9] Nalini Anantharaman and Mostafa Sabri. Quantum ergodicity : from spectral to spatial delocalization arxiv:1704.02766. *preprint*, 2017.
- [10] Nalini Anantharaman and Mostafa Sabri. Quantum ergodicity for the Anderson model on regular graphs. *J. Math. Phys.*, 58(9):091901, 10, 2017.
- [11] Nalini Anantharaman and Mostafa Sabri. Recent results of quantum ergodicity on graphs and further investigation. *preprint*, *arXiv:1711.07666*, 2017.
- [12] Agnes Backhausz and Balazs Szegedy. On the almost eigenvectors of random regular graphs arxiv:1607.04785. *preprint*, 2016.
- [13] Roland Bauerschmidt, Jiaoyang Huang, Antti Knowles, and Horng-Tzer Yau. Bulk eigenvalue statistics for random regular graphs arxiv:1505.06700. 2015.
- [14] Roland Bauerschmidt, Jiaoyang Huang, and Horng-Tzer Yau. Local kesten–mckay law for random regular graphs arxiv:1609.09052. 2016.
- [15] Roland Bauerschmidt, Antti Knowles, and Horng-Tzer Yau. Local semicircle law for random regular graphs. *Comm. Pure Appl. Math.*, 70(10):1898–1960, 2017.

- [16] Itai Benjamini and Oded Schramm. Recurrence of distributional limits of finite planar graphs. *Electron. J. Probab.*, 6:no. 23, 13, 2001.
- [17] Pierre H. Bérard. On the wave equation on a compact Riemannian manifold without conjugate points. *Math. Z.*, 155(3):249–276, 1977.
- [18] G. Berkolaiko, J. P. Keating, and U. Smilansky. Quantum ergodicity for graphs related to interval maps. *Comm. Math. Phys.*, 273(1):137–159, 2007.
- [19] G. Berkolaiko, J. P. Keating, and B. Winn. No quantum ergodicity for star graphs. *Comm. Math. Phys.*, 250(2):259–285, 2004.
- [20] Michael Berry. Regular and irregular semiclassical wavefunctions. *J. Phys. A*, 10(12):2083–2091, 1977.
- [21] Michael Berry and Michael Tabor. Level clustering in the regular spectrum. *Proc. Royal Soc. A*, pages 375–394, 1977.
- [22] Matthew Blair and Christopher Sogge. Concerning toponogov’s theorem and logarithmic improvement of estimates of eigenfunctions, to appear in *J. Diff. Geom.* 2015.
- [23] Matthew Blair and Christopher Sogge. Logarithmic improvements in l^p bounds for eigenfunctions at the critical exponent in the presence of nonpositive curvature. 2017.
- [24] Matthew D. Blair and Christopher D. Sogge. Refined and Microlocal Keakeya–Nikodym Bounds of Eigenfunctions in Higher Dimensions. *Comm. Math. Phys.*, 356(2):501–533, 2017.
- [25] Siegfried Böcherer, Peter Sarnak, and Rainer Schulze-Pillot. Arithmetic and equidistribution of measures on the sphere. *Comm. Math. Phys.*, 242(1-2):67–80, 2003.
- [26] O. Bohigas, M.-J. Giannoni, and C. Schmit. Characterization of chaotic quantum spectra and universality of level fluctuation laws. *Phys. Rev. Lett.*, 52(1):1–4, 1984.
- [27] Oriol Bohigas. Random matrix theories and chaotic dynamics. In *Chaos et physique quantique (Les Houches, 1989)*, pages 87–199. North-Holland, Amsterdam, 1991.
- [28] Niels Bohr. On the constitution of atoms and molecules. *Philosophical Magazine*, 26:1–24, 1913.
- [29] Jens Bolte and Rainer Glaser. A semiclassical Egorov theorem and quantum ergodicity for matrix valued operators. *Comm. Math. Phys.*, 247(2):391–419, 2004.
- [30] Yannick Bonthouneau. The Θ function and the Weyl law on manifolds without conjugate points. *Doc. Math.*, 22:1275–1283, 2017.
- [31] Yannick Bonthouneau and Steve Zelditch. Quantum ergodicity for Eisenstein functions. *C. R. Math. Acad. Sci. Paris*, 354(9):907–911, 2016.
- [32] David Borthwick. *Spectral theory of infinite-area hyperbolic surfaces*, volume 256 of *Progress in Mathematics*. Birkhäuser Boston, Inc., Boston, MA, 2007.
- [33] P. Bourgade and H.-T. Yau. The eigenvector moment flow and local quantum unique ergodicity. *preprint*, 2013.
- [34] Jean Bourgain and Elon Lindenstrauss. Entropy of quantum limits. *Comm. Math. Phys.*, 233(1):153–171, 2003.
- [35] Louis Brillouin. La mécanique ondulatoire de schrödinger; une méthode générale de résolution par approximations successives. *C. R. A. S.*, (183):24–26, 1926.
- [36] Shimon Brooks and Étienne Le Masson. l^p -norms of eigenfunctions on regular graphs and on the sphere arxiv:1710.10922 d. *preprint*, 2017.
- [37] Shimon Brooks, Étienne Le Masson, and Elon Lindenstrauss. Quantum ergodicity and averaging operators on the sphere. *Int. Math. Res. Not.*, 2015. To appear, arXiv:1505.03887.
- [38] Shimon Brooks and Elon Lindenstrauss. Non-localization of eigenfunctions on large regular graphs. *Israel J. Math.*, 193(1):1–14, 2013.
- [39] Shimon Brooks and Elon Lindenstrauss. Joint quasimodes, positive entropy, and quantum unique ergodicity. *Invent. Math.*, 198(1):219–259, 2014.
- [40] Y. Colin de Verdière. Ergodicité et fonctions propres du laplacien. *Comm. Math. Phys.*, 102(3):497–502, 1985.

- [41] Y. Colin de Verdière. Semi-classical measures on quantum graphs and the gauss map of the determinant manifold. *Ann. H. Poincaré*, to appear, 2013.
- [42] Yves Colin de Verdière, Luc Hillairet, and Emmanuel Trélat. Spectral asymptotics for sub-riemannian laplacians. i: quantum ergodicity and quantum limits in the 3d contact case. *preprint*, 2015.
- [43] Dominique Delande. Chaos in atomic and molecular physics. In *Chaos et physique quantique (Les Houches, 1989)*, pages 665–726. North-Holland, Amsterdam, 1991.
- [44] Persi Diaconis and Daniel Stroock. Geometric bounds for eigenvalues of Markov chains. *Ann. Appl. Probab.*, 1(1):36–61, 1991.
- [45] Ioana Dumitriu and Soumik Pal. Sparse regular random graphs: spectral density and eigenvectors. *Ann. Probab.*, 40(5):2197–2235, 2012.
- [46] Semyon Dyatlov and Colin Guillarmou. Microlocal limits of plane waves and Eisenstein functions. *Ann. Sci. Éc. Norm. Supér. (4)*, 47(2):371–448, 2014.
- [47] Semyon Dyatlov and Long Jin. Semiclassical measures on hyperbolic surfaces have full support [arxiv:1705.05019](https://arxiv.org/abs/1705.05019). *preprint*, 2017.
- [48] Manfred Einsiedler and Elon Lindenstrauss. Diagonalizable flows on locally homogeneous spaces and number theory. In *International Congress of Mathematicians. Vol. II*, pages 1731–1759. Eur. Math. Soc., Zürich, 2006.
- [49] Albert Einstein. über einen die erzeugung und verwandlung des lichts betre enden heuristischen gesichtspunkt. *Annalen der Physik*, 17, 1905.
- [50] Albert Einstein. Zum quantensatz von sommerfeld und epstein. *Verhandl. deut. physik. Ges.*, 1917.
- [51] László Erdős, Antti Knowles, Horng-Tzer Yau, and Jun Yin. Spectral statistics of Erdős-Rényi graphs I: Local semicircle law. *Ann. Probab.*, 41(3B):2279–2375, 2013.
- [52] László Erdős, Benjamin Schlein, and Horng-Tzer Yau. Local semicircle law and complete delocalization for Wigner random matrices. *Comm. Math. Phys.*, 287(2):641–655, 2009.
- [53] László Erdős, Benjamin Schlein, and Horng-Tzer Yau. Semicircle law on short scales and delocalization of eigenvectors for Wigner random matrices. *Ann. Probab.*, 37(3):815–852, 2009.
- [54] Frédéric Faure and Stéphane Nonnenmacher. On the maximal scarring for quantum cat map eigenstates. *Comm. Math. Phys.*, 245(1):201–214, 2004.
- [55] Frédéric Faure, Stéphane Nonnenmacher, and Stephan De Bièvre. Scarred eigenstates for quantum cat maps of minimal periods. *Comm. Math. Phys.*, 239(3):449–492, 2003.
- [56] Leander Geisinger. Convergence of the density of states and delocalization of eigenvectors on random regular graphs. *J. Spectr. Theory*, 5(4):783–827, 2015.
- [57] Patrick Gérard and Éric Leichtnam. Ergodic properties of eigenfunctions for the Dirichlet problem. *Duke Math. J.*, 71(2):559–607, 1993.
- [58] S. Gnuzmann, J. P. Keating, and F. Pietet. Eigenfunction statistics on quantum graphs. *Ann. Physics*, 325(12):2595–2640, 2010.
- [59] Colin Guillarmou and Frédéric Naud. Equidistribution of Eisenstein series on convex co-compact hyperbolic manifolds. *Amer. J. Math.*, 136(2):445–479, 2014.
- [60] Xiaolong Han. Small scale quantum ergodicity in negatively curved manifolds. *Nonlinearity*, 28(9):3263–3288, 2015.
- [61] Xiaolong Han. Small scale equidistribution of random eigenbases. *Comm. Math. Phys.*, 349(1):425–440, 2017.
- [62] Andrew Hassell. Ergodic billiards that are not quantum unique ergodic. *Ann. of Math. (2)*, 171(1):605–619, 2010. With an appendix by the author and Luc Hillairet.
- [63] Andrew Hassell and Melissa Tacy. Improvement of eigenfunction estimates on manifolds of nonpositive curvature. *Forum Math.*, 27(3):1435–1451, 2015.
- [64] Werner Heisenberg. über quantentheoretische umdeutung kinematischer und mechanischer beziehungen. *Zeitschrift f. Physik*, 33:879–893, 1925.

- [65] B. Helffer, A. Martinez, and D. Robert. Ergodicité et limite semi-classique. *Comm. Math. Phys.*, 109(2):313–326, 1987.
- [66] Eric Heller. Wavepacket dynamics and quantum chaos. In *Chaos et physique quantique (Les Houches, 1989)*, pages 549–661. North-Holland, Amsterdam, 1991.
- [67] Hamid Hezari and Gabriel Rivière. L^p norms, nodal sets, and quantum ergodicity. *Adv. Math.*, 290:938–966, 2016.
- [68] Maxime Ingremeau. Distorted plane waves on manifolds of nonpositive curvature. *Comm. Math. Phys.*, 350(2):845–891, 2017.
- [69] Dmitry Jakobson. Quantum unique ergodicity for Eisenstein series on $\mathrm{PSL}_2(\mathbf{Z}) \backslash \mathrm{PSL}_2(\mathbf{R})$. *Ann. Inst. Fourier (Grenoble)*, 44(5):1477–1504, 1994.
- [70] Dmitry Jakobson, Yuri Safarov, and Alexander Strohmaier. The semiclassical theory of discontinuous systems and ray-splitting billiards. *Amer. J. Math.*, 137(4):859–906, 2015. With an appendix by Yves Colin de Verdière.
- [71] Dmitry Jakobson and Alexander Strohmaier. High energy limits of Laplace-type and Dirac-type eigenfunctions and frame flows. *Comm. Math. Phys.*, 270(3):813–833, 2007.
- [72] Dmitry Jakobson, Alexander Strohmaier, and Steve Zelditch. On the spectrum of geometric operators on Kähler manifolds. *J. Mod. Dyn.*, 2(4):701–718, 2008.
- [73] J. P. Keating. Quantum graphs and quantum chaos. In *Analysis on graphs and its applications*, volume 77 of *Proc. Sympos. Pure Math.*, pages 279–290. Amer. Math. Soc., Providence, RI, 2008.
- [74] J. P. Keating, J. Marklof, and B. Winn. Value distribution of the eigenfunctions and spectral determinants of quantum star graphs. *Comm. Math. Phys.*, 241(2-3):421–452, 2003.
- [75] Harry Kesten. Symmetric random walks on groups. *Trans. Amer. Math. Soc.*, 92:336–354, 1959.
- [76] Abel Klein. Extended states in the Anderson model on the Bethe lattice. *Adv. Math.*, 133(1):163–184, 1998.
- [77] Tsampikos Kottos and Uzy Smilansky. Quantum chaos on graphs. *Phys. Rev. Lett.*, 79:4794–7, 1997.
- [78] Tsampikos Kottos and Uzy Smilansky. Periodic orbit theory and spectral statistics for quantum graphs. *Ann. Physics*, 274(1):76–124, 1999.
- [79] Hendrik Anthony Kramers. Wellenmechanik und halbzahlige quantisierung. *Zeitschrift f. Physik*, (39):828–840, 1926.
- [80] Étienne Le Masson and Tuomas Sahlsten. Quantum ergodicity and benjamini-schramm convergence of hyperbolic surfaces. [arxiv:1605.05720](https://arxiv.org/abs/1605.05720). *Duke Math. J.*, to appear.
- [81] F. Ledrappier and L.-S. Young. The metric entropy of diffeomorphisms. I. Characterization of measures satisfying Pesin’s entropy formula. *Ann. of Math. (2)*, 122(3):509–539, 1985.
- [82] F. Ledrappier and L.-S. Young. The metric entropy of diffeomorphisms. II. Relations between entropy, exponents and dimension. *Ann. of Math. (2)*, 122(3):540–574, 1985.
- [83] Stephen Lester and Zeév Rudnick. Small scale equidistribution of eigenfunctions on the torus. *Comm. Math. Phys.*, 350(1):279–300, 2017.
- [84] Elon Lindenstrauss. Invariant measures and arithmetic quantum unique ergodicity. *Ann. of Math. (2)*, 163(1):165–219, 2006.
- [85] Elon Lindenstrauss. Equidistribution in homogeneous spaces and number theory. In *Proceedings of the International Congress of Mathematicians. Volume I*, pages 531–557. Hindustan Book Agency, New Delhi, 2010.
- [86] A. Lubotzky, R. Phillips, and P. Sarnak. Ramanujan graphs. *Combinatorica*, 8(3):261–277, 1988.
- [87] W. Luo and P. Sarnak. Number variance for arithmetic hyperbolic surfaces. *Comm. Math. Phys.*, 161(2):419–432, 1994.
- [88] Jens Marklof. Pair correlation densities of inhomogeneous quadratic forms. *Ann. of Math. (2)*, 158(2):419–471, 2003.
- [89] Brendan D. McKay. The expected eigenvalue distribution of a large regular graph. *Linear Algebra Appl.*, 40:203–216, 1981.

- [90] Brendan D. McKay, Nicholas C. Wormald, and Beata Wysocka. Short cycles in random regular graphs. *Electron. J. Combin.*, 11(1):Research Paper 66, 12 pp. (electronic), 2004.
- [91] Maryam Mirzakhani. On Weil-Petersson volumes and geometry of random hyperbolic surfaces. In *Proceedings of the International Congress of Mathematicians. Volume II*, pages 1126–1145. Hindustan Book Agency, New Delhi, 2010.
- [92] Maryam Mirzakhani. Growth of Weil-Petersson volumes and random hyperbolic surfaces of large genus. *J. Differential Geom.*, 94(2):267–300, 2013.
- [93] Max Planck. über eine verbesserung der wienschen spektralgleichung. *Verhandl. deut. physik. Ges.*, 13:202–204, 1900.
- [94] Gabriel Rivière. Entropy of semiclassical measures in dimension 2. *Duke. Math. J.*, 155(2):271–335, 2010.
- [95] Zeév Rudnick and Peter Sarnak. The behaviour of eigenstates of arithmetic hyperbolic manifolds. *Comm. Math. Phys.*, 161(1):195–213, 1994.
- [96] Peter Sarnak. Arithmetic quantum chaos. In *The Schur lectures (1992) (Tel Aviv)*, volume 8 of *Israel Math. Conf. Proc.*, pages 183–236. Bar-Ilan Univ., Ramat Gan, 1995.
- [97] Peter Sarnak. Values at integers of binary quadratic forms. In *Harmonic analysis and number theory (Montreal, PQ, 1996)*, volume 21 of *CMS Conf. Proc.*, pages 181–203. Amer. Math. Soc., Providence, RI, 1997.
- [98] Peter Sarnak. Spectra of hyperbolic surfaces. *Bull. Amer. Math. Soc. (N.S.)*, 40(4):441–478, 2003.
- [99] Erwin Schrödinger. Quantisierung als eigenwertproblem (erste mitteilung). *Annalen der Physik (4)*, 79:361–376, 1926.
- [100] Erwin Schrödinger. Quantisierung als eigenwertproblem (zweite mitteilung). *Annalen der Physik (4)*, 79:489–527, 1926.
- [101] Erwin Schrödinger. über das verhältnis der heisenberg–born–jordanschen quantenmechanik zu der meinen. *Annalen der Physik (4)*, 79:734–756, 1926.
- [102] Atle Selberg. Remarks on the distribution of poles of Eisenstein series. In *Festschrift in honor of I. I. Piatetski-Shapiro on the occasion of his sixtieth birthday, Part II (Ramat Aviv, 1989)*, volume 3 of *Israel Math. Conf. Proc.*, pages 251–278. Weizmann, Jerusalem, 1990.
- [103] Martin Sieber and Klaus Richter. Correlations between periodic orbits and their rôle in spectral statistics. *Physica Scripta*, T90, 2001.
- [104] Uzy Smilansky. Quantum chaos on discrete graphs. *J. Phys. A*, 40(27):F621–F630, 2007.
- [105] Uzy Smilansky. Discrete graphs – a paradigm model for quantum chaos. *Séminaire Poincaré*, XIV:1–26, 2010.
- [106] A. I. Šnirel'man. Ergodic properties of eigenfunctions. *Uspehi Mat. Nauk*, 29(6(180)):181–182, 1974.
- [107] Christopher D. Sogge. Concerning the L^p norm of spectral clusters for second-order elliptic operators on compact manifolds. *J. Funct. Anal.*, 77(1):123–138, 1988.
- [108] Christopher D. Sogge and Steve Zelditch. Riemannian manifolds with maximal eigenfunction growth. *Duke Math. J.*, 114(3):387–437, 2002.
- [109] Christopher D. Sogge and Steve Zelditch. Focal points and sup-norms of eigenfunctions. *Rev. Mat. Iberoam.*, 32(3):971–994, 2016.
- [110] Christopher D. Sogge and Steve Zelditch. Focal points and sup-norms of eigenfunctions II: the two-dimensional case. *Rev. Mat. Iberoam.*, 32(3):995–999, 2016.
- [111] K. Soundararajan. Quantum unique ergodicity and number theory. In *Proceedings of the International Congress of Mathematicians. Volume II*, pages 357–382. Hindustan Book Agency, New Delhi, 2010.
- [112] Kannan Soundararajan. Quantum unique ergodicity for $SL_2(\mathbb{Z}) \backslash \mathbb{H}$. *Ann. of Math. (2)*, 172(2):1529–1538, 2010.
- [113] Linh V. Tran, Van H. Vu, and Ke Wang. Sparse random graphs: eigenvalues and eigenvectors. *Random Structures Algorithms*, 42(1):110–134, 2013.
- [114] Jeffrey M. VanderKam. L^∞ norms and quantum ergodicity on the sphere. *Internat. Math. Res. Notices*, (7):329–347, 1997.

- [115] Gregor Wentzel. Eine verallgemeinerung der quantenbedingungen für die zwecke der wellenmechanik. *Zeitschrift f. Physik*, (38):518–529, 1926.
- [116] Steven Zelditch. Uniform distribution of eigenfunctions on compact hyperbolic surfaces. *Duke Math. J.*, 55(4):919–941, 1987.
- [117] Steven Zelditch. Mean Lindelöf hypothesis and equidistribution of cusp forms and Eisenstein series. *J. Funct. Anal.*, 97(1):1–49, 1991.
- [118] Steven Zelditch. Quantum ergodicity on the sphere. *Comm. Math. Phys.*, 146(1):61–71, 1992.
- [119] Steven Zelditch and Maciej Zworski. Ergodicity of eigenfunctions for ergodic billiards. *Comm. Math. Phys.*, 175(3):673–682, 1996.

IRMA, 7 rue René-Descartes, 67084 STRASBOURG CEDEX, FRANCE

E-mail address: anantharaman@unistra.fr